

# Emerging phases and phase transitions in quantum matter

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- Condensed matter physics: complexity and emerging phenomena
  - Phase transitions and quantum phase transitions
  - Emerging phases close to quantum critical points

# Acknowledgements



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# What is condensed matter physics?

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**Condensed Matter Physics** (Wikipedia):

*deals with the **macroscopic** properties of matter; in particular ... the “condensed” phases that appear whenever the number of constituents in a system is **large** and their interactions ... are **strong***

**Traditionally: Physics of solids and liquids**

- What is the structure of crystals?
- How do solids melt or liquids evaporate?
- Why do some materials conduct an electric current and others do not?

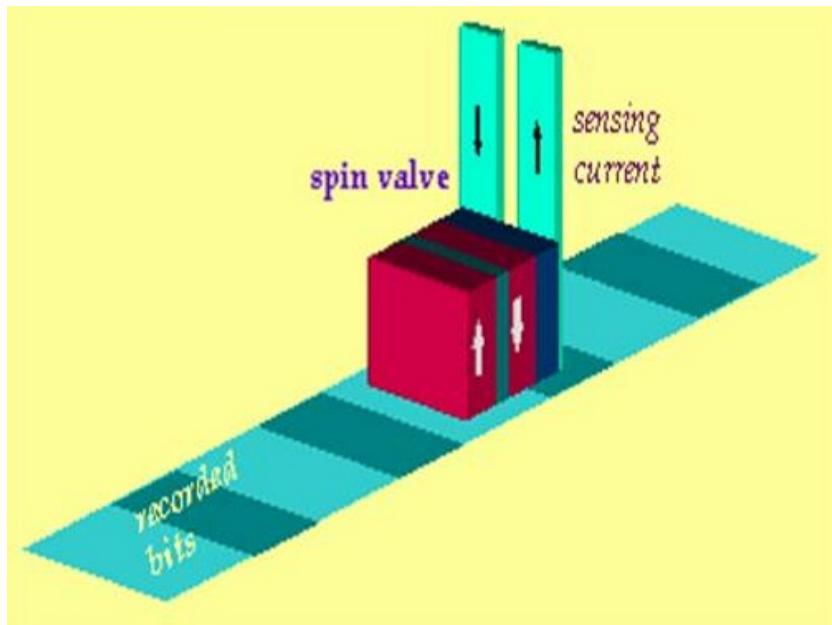
**Today: all systems consisting of a large number of interacting constituents**

- biological systems: biomolecules, DNA, membranes, cells
- geological systems: earthquakes
- economical systems: fluctuations of stock markets, currencies

# Why condensed matter physics?

## Applications: "Helps you to make stuff."

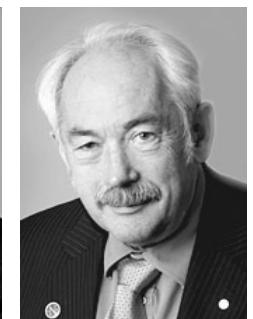
- semiconductors, transistors, microchips
- magnetic recording devices
- liquid crystal displays
- to come: qbits for quantum computers



Read head, based on **Giant Magnetoresistance effect**  
(A. Fert + P. Grünberg, Physics Nobel Prize 2007)



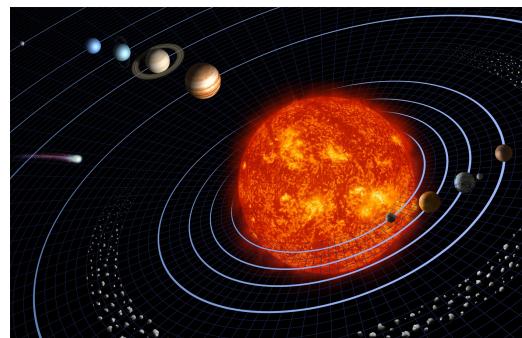
Maglev train using levitation by superconducting magnets, can go faster than 350 mph



# Why condensed matter physics II

Directions of fundamental physics research :

**Astrophysics and cosmology:**  
increasing length and time scales  
*“physics of the very large”*



**Atomic, nuclear and elementary particle physics:**  
decreasing length and time scales  
*“physics of the very small”*

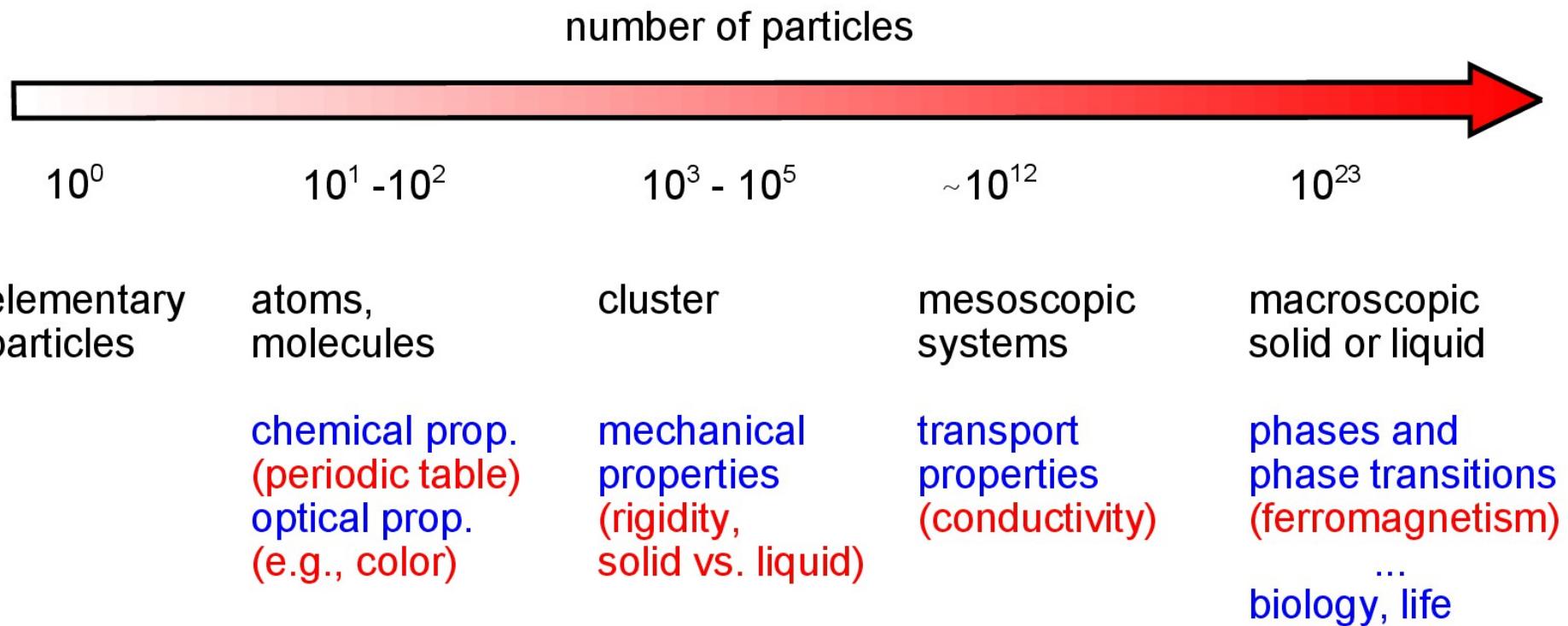


Particle accelerator at Fermilab

What fundamental direction does condensed matter research explore?  
*“physics of the very complex”*

# Emerging phenomena and the axis of complexity

*“More is different!”*



## Emerging phenomena:

When **large numbers** of particles strongly **interact**, qualitatively new properties of matter **emerge** at every level of complexity

# New states of quantum matter and where to find them

## Quantum matter:

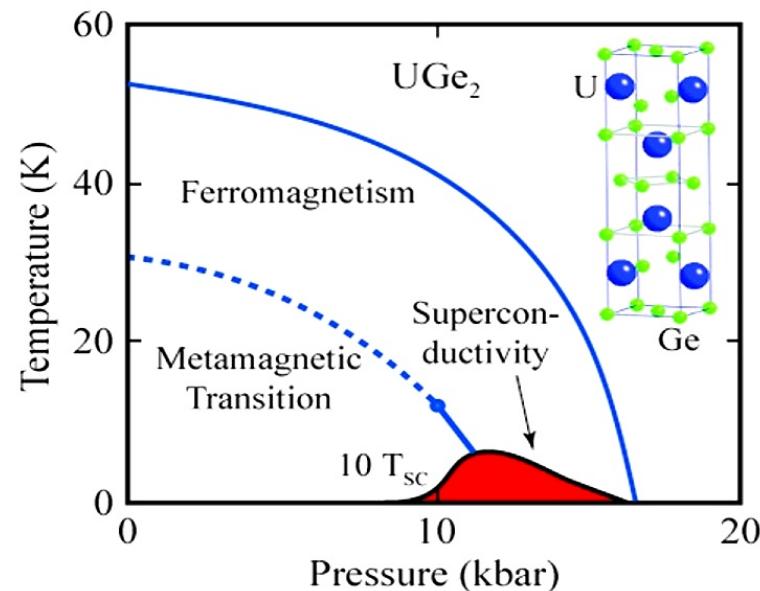
- matter with emerging **macroscopic** properties that are intrinsically **quantum** (superconductors, superfluids, fractional quantum Hall states, spin liquids)



### at low temperatures

$$F = E - TS$$

- thermal motion is suppressed
- new types of order can form



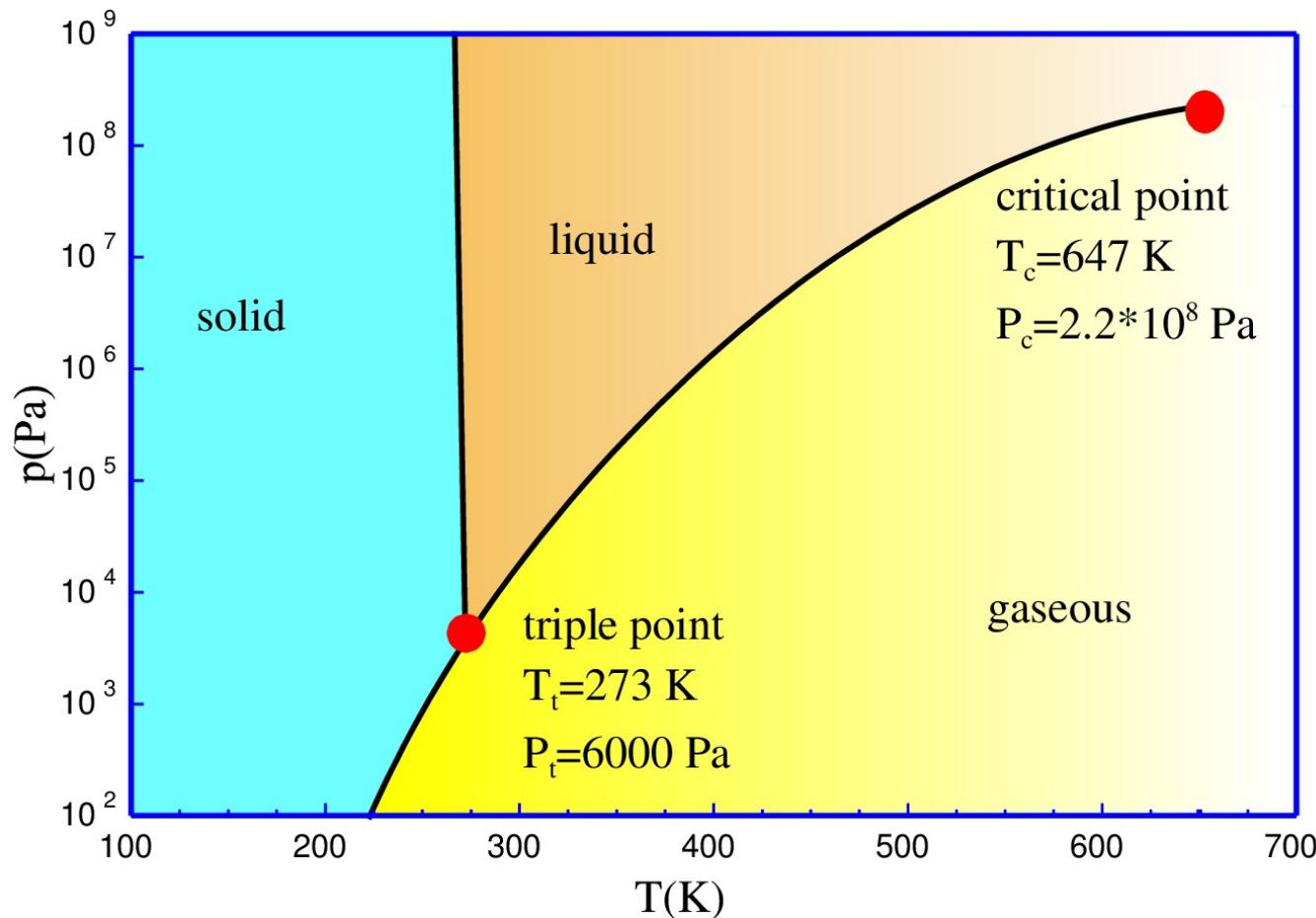
### at boundaries of existing phases

- two types of order compete, suppress each other
- novel type of order may appear

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- Condensed matter physics: complexity and emerging phenomena
- **Phase transitions and quantum phase transitions**
  - Novel phases close to quantum critical points

# Phase diagram of water



## Phase transition:

singularity in thermodynamic quantities as functions of external parameters

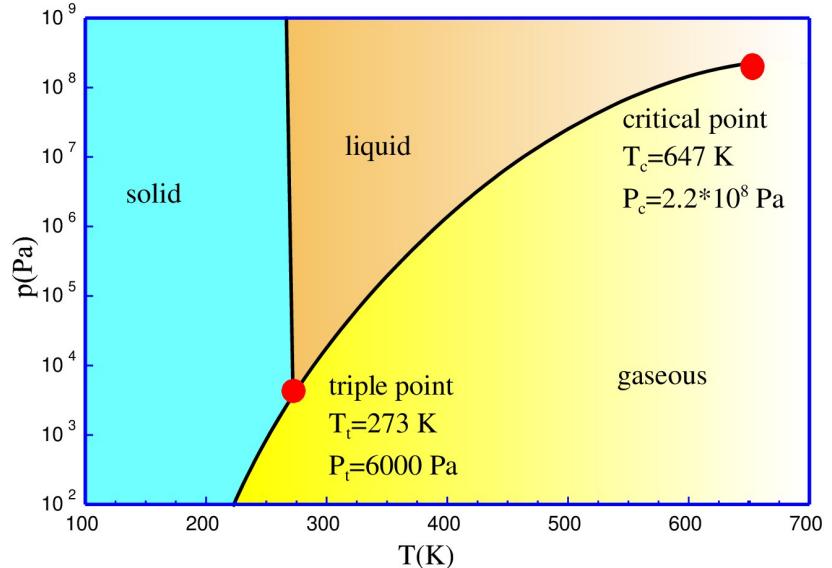
# Phase transitions: 1st order vs. continuous

## 1st order phase transition:

phase coexistence, latent heat,  
short range spatial and time correlations

## Continuous transition (critical point):

no phase coexistence, no latent heat,  
infinite range correlations of fluctuations



## Critical behavior at continuous transitions:

diverging correlation length  $\xi \sim |T - T_c|^{-\nu}$  and time  $\xi_\tau \sim \xi^z \sim |T - T_c|^{-\nu z}$

- Manifestation: critical opalescence (Andrews 1869)

**Universality: critical exponents are independent of microscopic details**

# Critical opalescence

## Binary liquid system:

e.g. hexane and methanol

$T > T_c \approx 36^\circ\text{C}$ : fluids are miscible

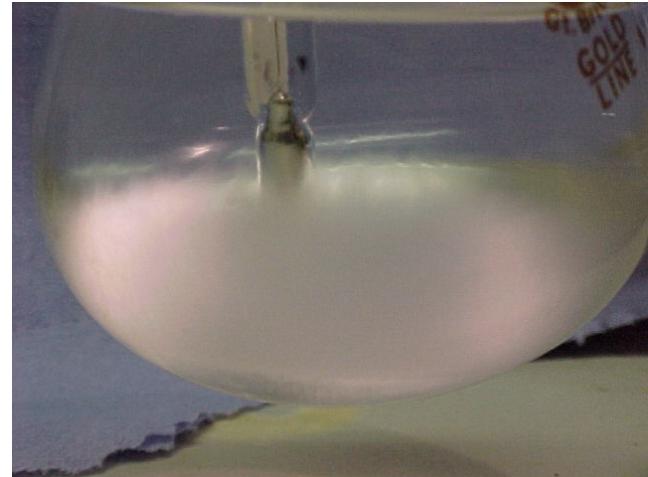
$T < T_c$ : fluids separate into two phases

$T \rightarrow T_c$ : length scale  $\xi$  of fluctuations grows

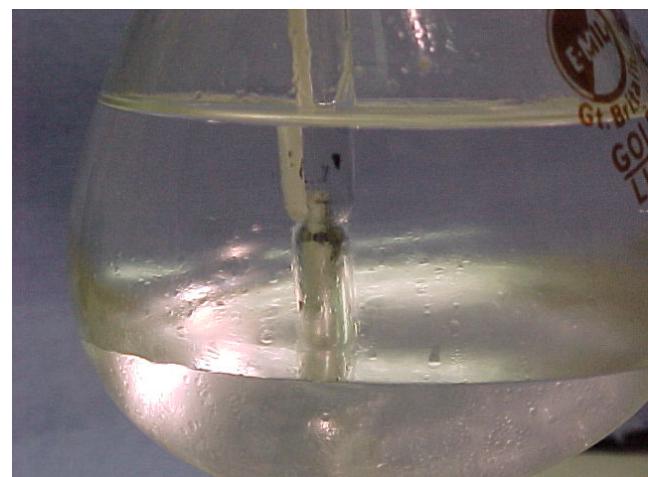
When  $\xi$  reaches the scale of a fraction of a micron (wavelength of light):

**strong light scattering**  
**fluid appears milky**

46°C



39°C



18°C



# How important is quantum mechanics close to a critical point?

## Two types of fluctuations:

thermal fluctuations (**thermal motion**), energy scale  $k_B T$

quantum fluctuations (**quantum zero-point motion**), energy scale  $\hbar\omega_c$

Quantum effects **unimportant** if  $\hbar\omega_c \ll k_B T$ .

## Critical slowing down:

$$\omega_c \sim 1/\xi_\tau \sim |T - T_c|^{\nu z} \rightarrow 0 \text{ at the critical point}$$

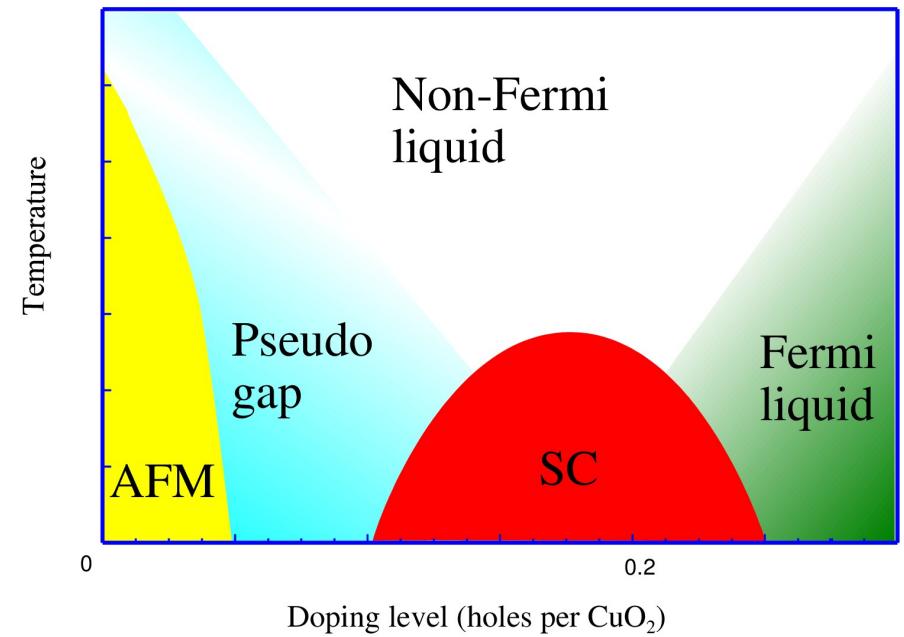
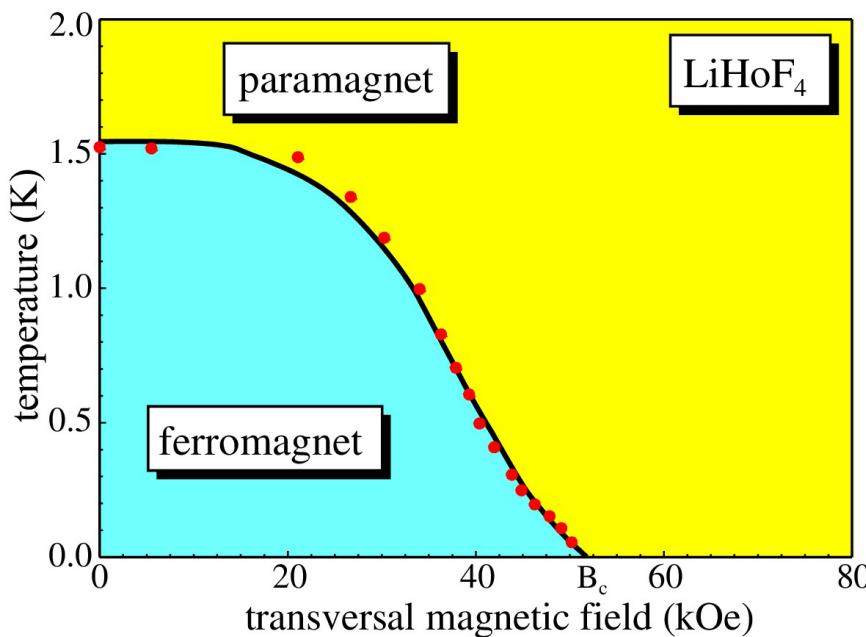
- ⇒ For any **nonzero** temperature, quantum fluctuations do **not** play a role close to the critical point
- ⇒ Quantum fluctuations **do** play a role at **zero** temperature

Zero-temperature continuous phase transitions constitute a special class of phase transitions, they are intrinsically quantum in nature

# Quantum phase transitions

occur at **zero temperature** as function of pressure, magnetic field, chemical composition, ...

driven by **quantum zero-point motion** rather than thermal fluctuations

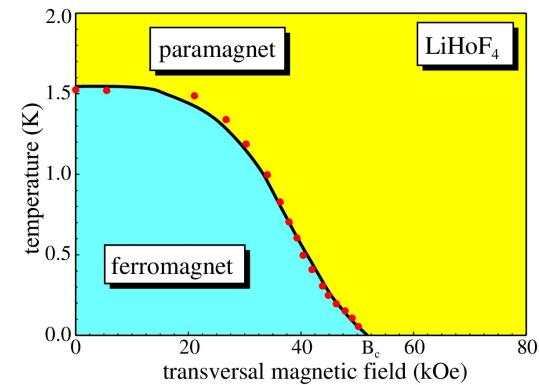


Phase diagrams of  $\text{LiHoF}_4$  and a typical high- $T_c$  superconductor such as  $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$

## Toy model: transverse field Ising model

Quantum spins  $S_i$  on a lattice: (c.f.  $\text{LiHoF}_4$ )

$$\begin{aligned} H &= -J \sum_i \mathbf{S}_i^z \mathbf{S}_{i+1}^z - h \sum_i \mathbf{S}_i^x \\ &= -J \sum_i \mathbf{S}_i^z \mathbf{S}_{i+1}^z - \frac{h}{2} \sum_i (\mathbf{S}_i^+ + \mathbf{S}_i^-) \end{aligned}$$



$J$ : exchange energy, favors parallel spins, i.e., ferromagnetic state

$h$ : transverse magnetic field, induces quantum fluctuations between up and down states, favors paramagnetic state

Limiting cases:

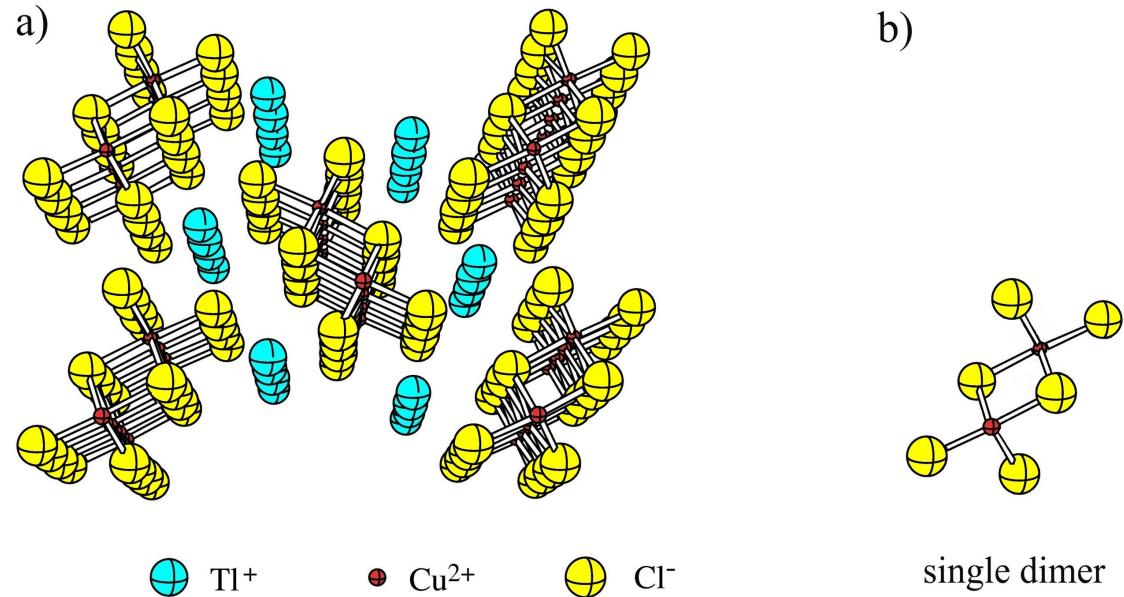
$|J| \gg |h|$  ferromagnetic ground state as in classical Ising magnet

$|J| \ll |h|$  paramagnetic ground state as for independent spins in a field

⇒ Quantum phase transition at  $|J| \sim |h|$  (in 1D, transition is at  $|J| = |h|$ )

# Magnetic quantum critical points of $\text{TiCuCl}_3$

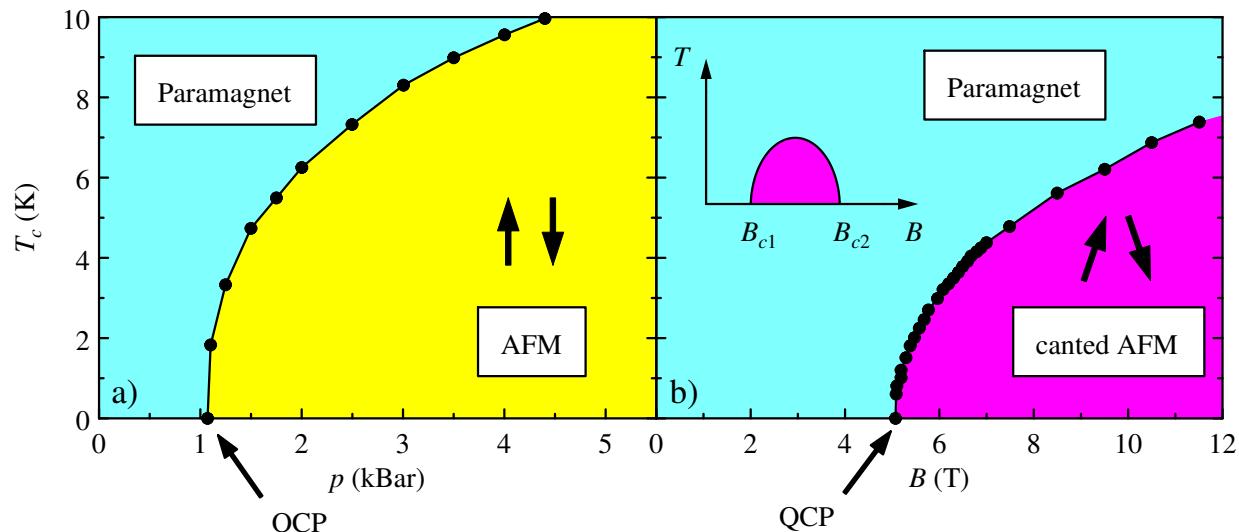
- $\text{TiCuCl}_3$  is magnetic insulator
- planar  $\text{Cu}_2\text{Cl}_6$  dimers form infinite double chains
- $\text{Cu}^{2+}$  ions carry spin-1/2 moment



## antiferromagnetic order

can be induced by

- applying pressure
- applying a magnetic field



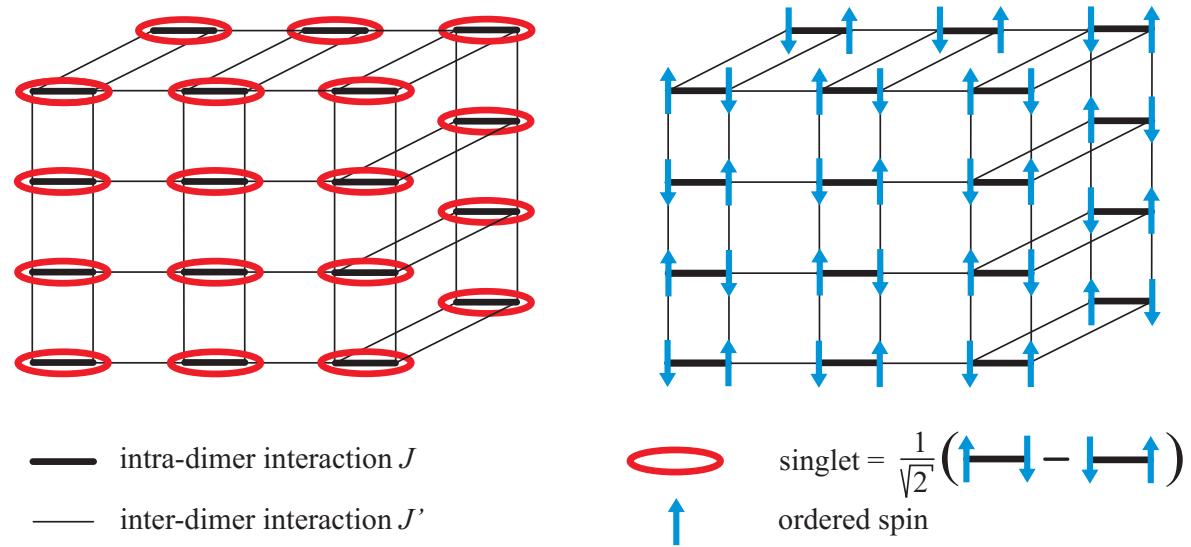
# Pressure-driven quantum phase transition in $\text{TlCuCl}_3$

## quantum Heisenberg model

$$H = \sum_{\langle ij \rangle} J_{ij} \vec{S}_i \cdot \vec{S}_j - \vec{h} \cdot \sum_i \vec{S}_i$$

$$J_{ij} = \begin{cases} J & \text{intra-dimer} \\ J' & \text{between dimers} \end{cases}$$

pressure changes ratio  $J/J'$



## Limiting cases:

$|J| \gg |J'|$  spins on each dimer form singlet  $\Rightarrow$  no magnetic order

low-energy excitations are “triplons” (single dimers in the triplet state)

$|J| \approx |J'|$  long-range antiferromagnetic order (Néel order)

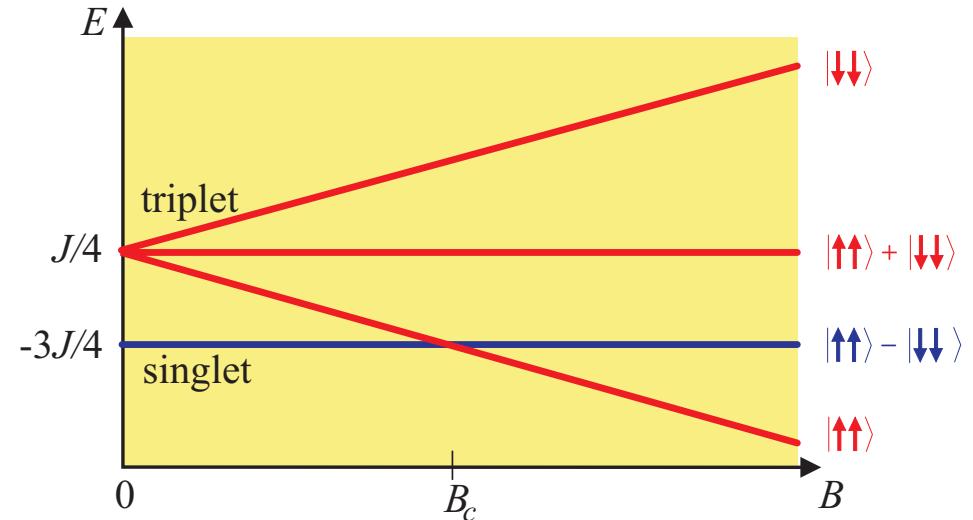
low-energy excitations are long-wavelength spin waves

⇒ quantum phase transition at some critical value of the ratio  $J/J'$

# Field-driven quantum phase transition in $\text{TiCuCl}_3$

## Single dimer in field:

- field does not affect singlet ground state but splits the triplet states
- ground state: singlet for  $B < B_c$  and (fully polarized) triplet for  $B > B_c$

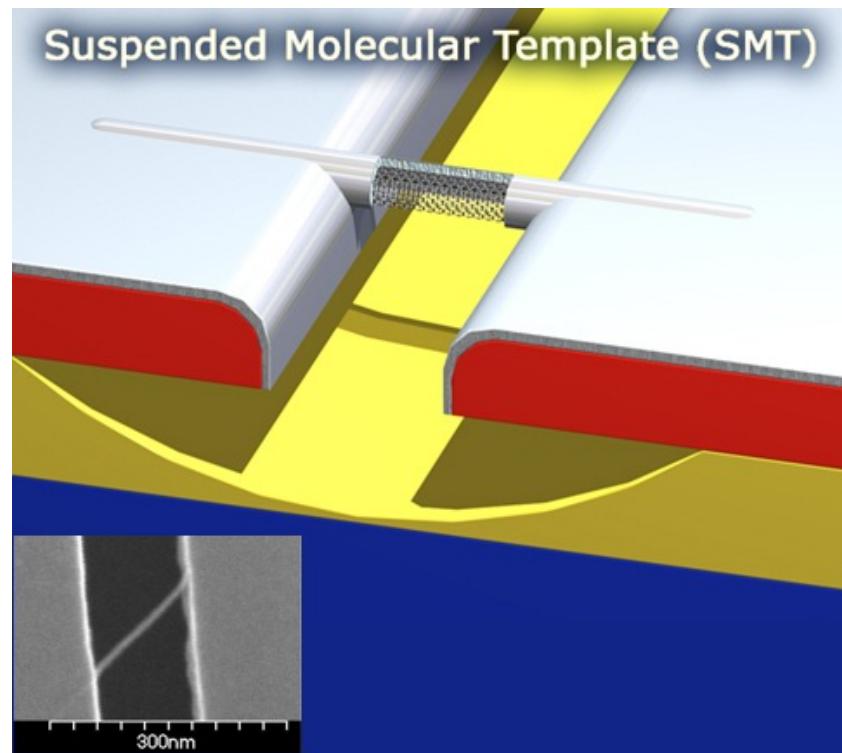
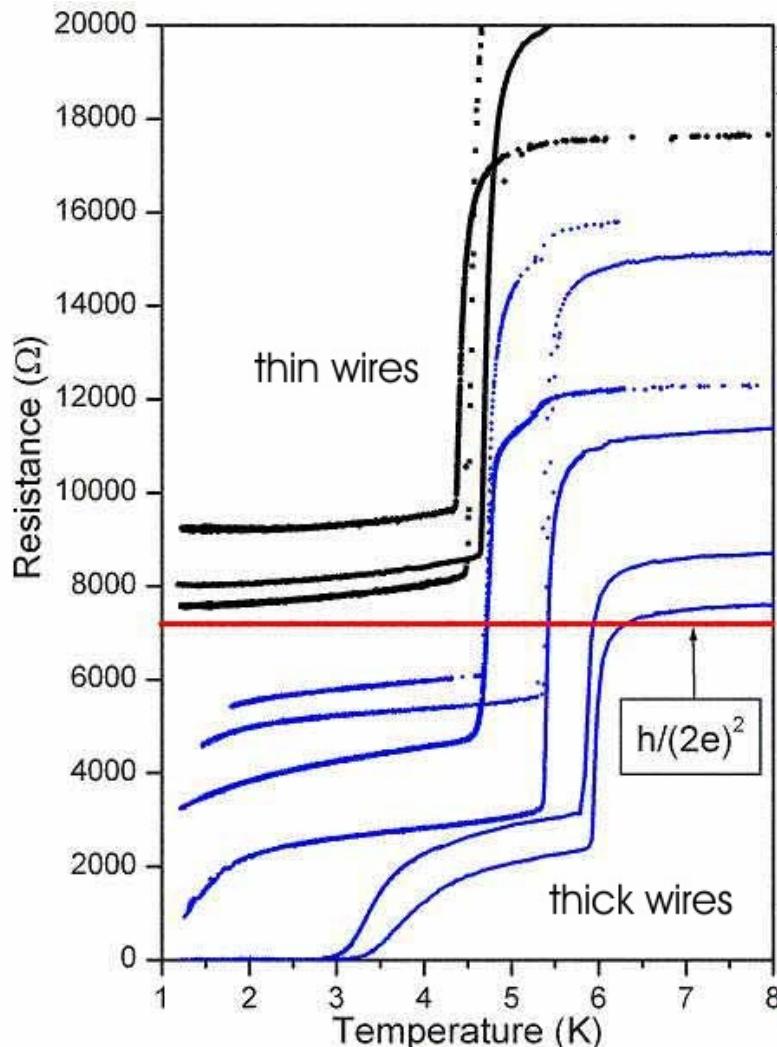


## Full Hamiltonian:

- singlet-triplet transition of isolated dimer splits into two transitions
- at  $B_{c1}$ , triplon gap closes, system is driven into ordered state (uniform magnetization  $\parallel$  to field and antiferromagnetic order  $\perp$  to field)
- “canted” antiferromagnet is Bose-Einstein condensate of triplons
- at  $B_{c2}$  system enters fully polarized state

# Superconductor-metal QPT in ultrathin nanowires

- ultrathin MoGe wires (width  $\sim 10$  nm)
- produced by molecular templating using a single carbon nanotube

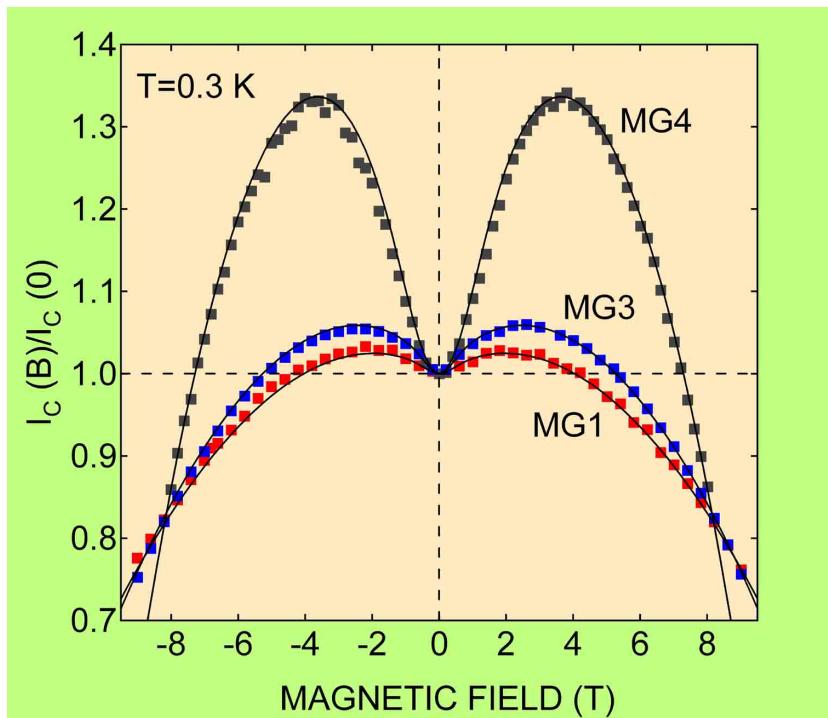


- thicker wires are superconducting at low temperatures
- thinner wires remain metallic

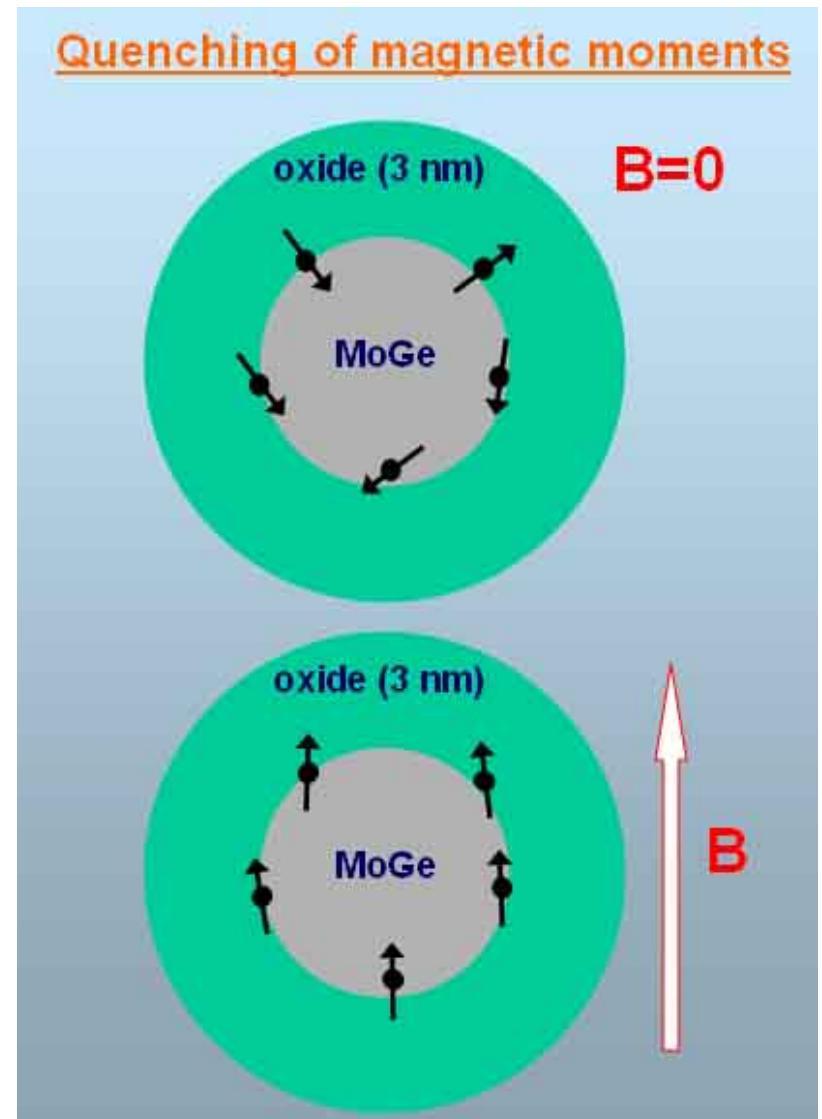
superconductor-metal QPT as function of wire thickness

# Pairbreaking mechanism

- pair breaking by surface magnetic impurities
- random impurity positions  
⇒ quenched **disorder**
- gapless excitations in metal phase  
⇒ Ohmic **dissipation**

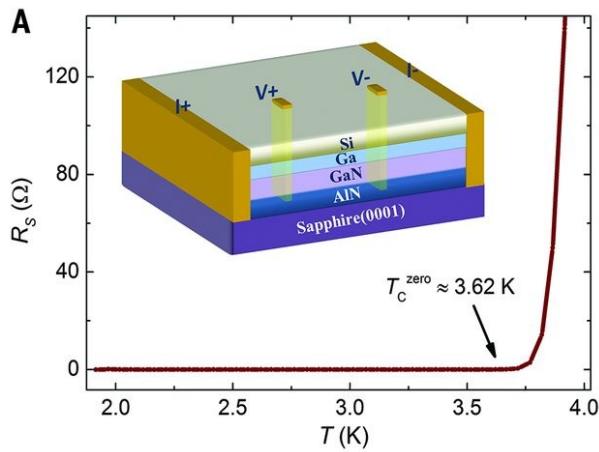


weak field enhances superconductivity



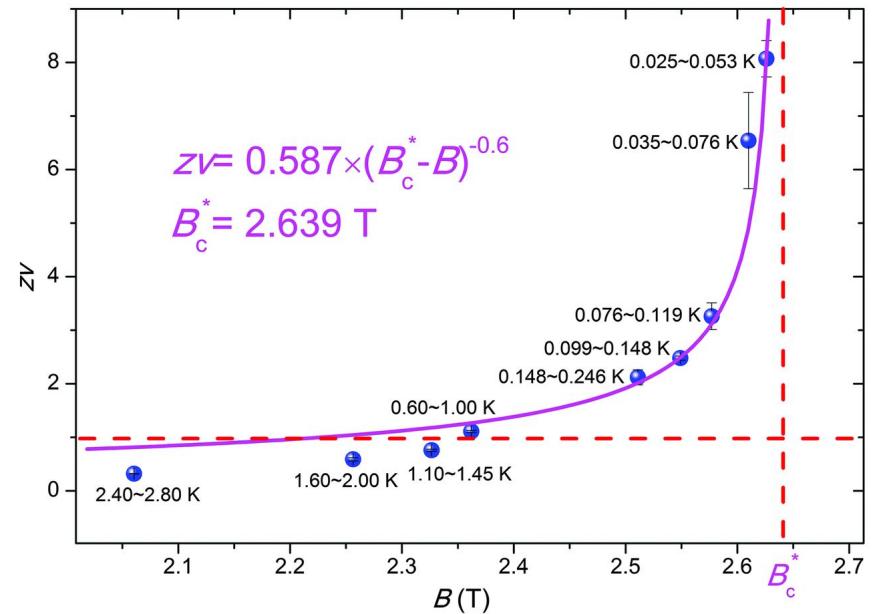
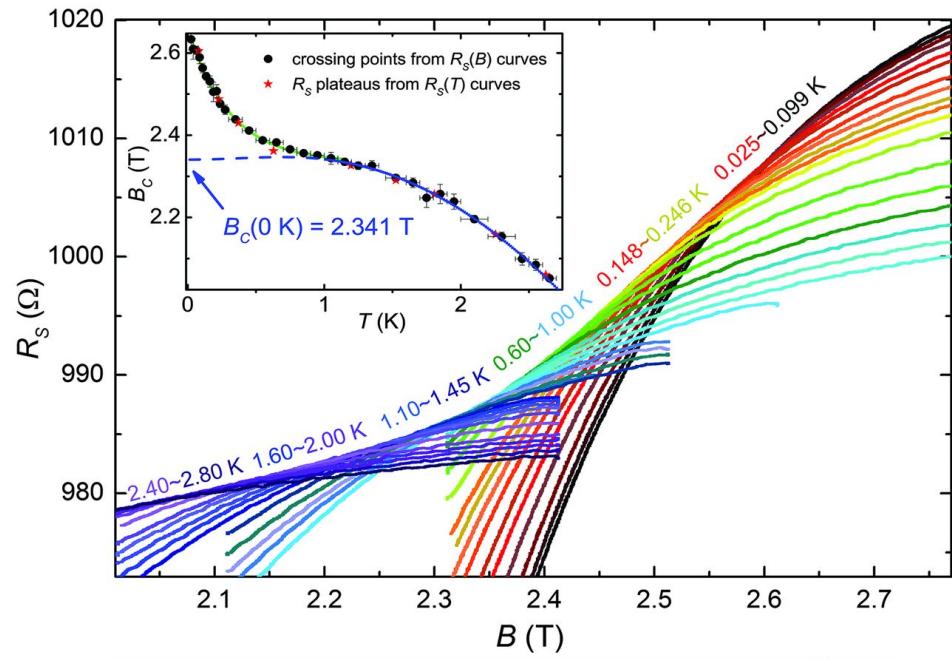
magnetic field aligns the impurities and reduces magnetic scattering

# Experiment: Ga thin films

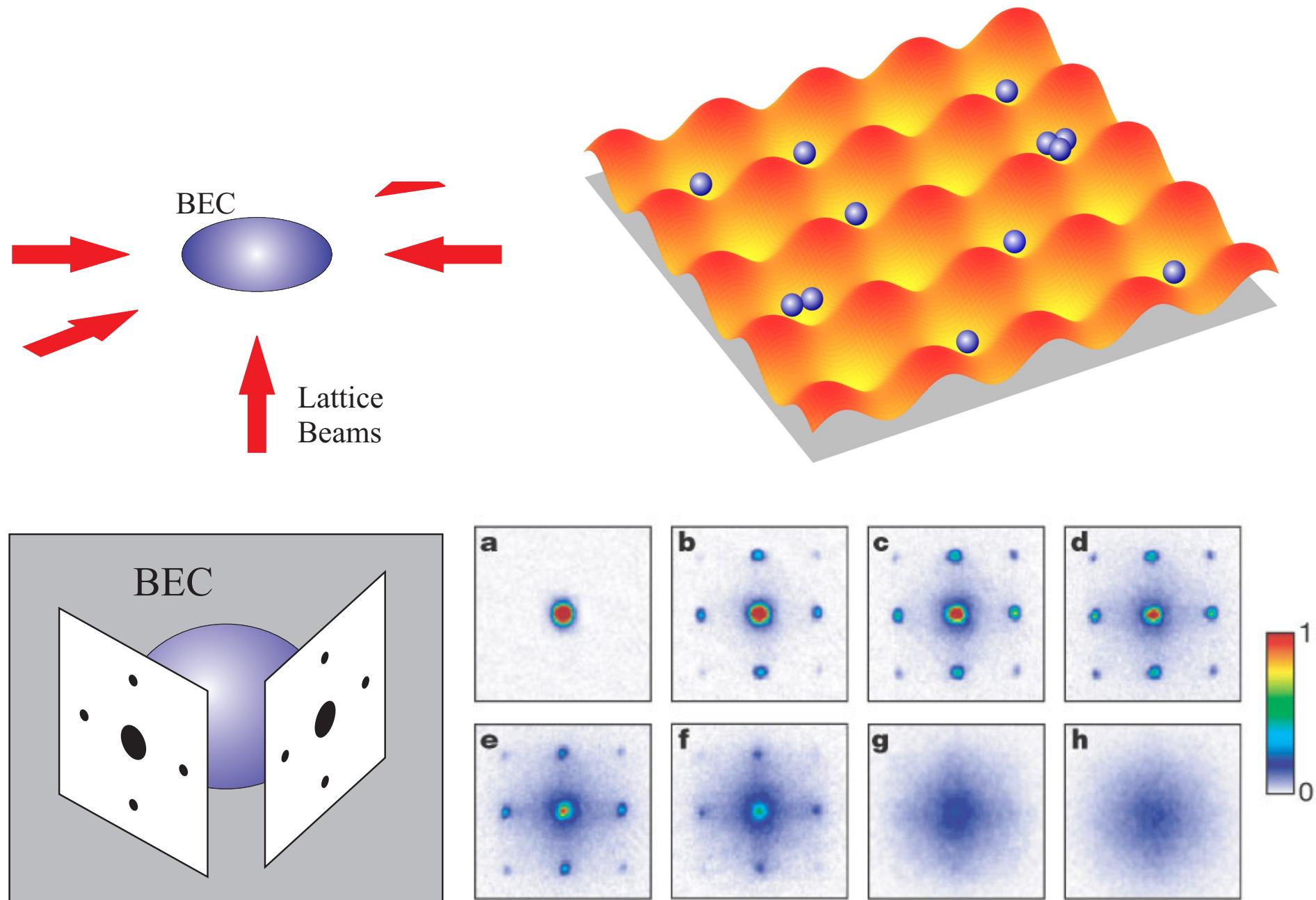


Xing et al., Science 350, 542 (2015)

- three-monolayer Ga films
- superconductivity below  $T_c \approx 3.62 \text{ K}$ , suppressed by magnetic field
- field-driven QPT well described by **2D infinite-randomness critical point**
- dynamical exponent **diverges** as  $z \sim |B - B_c|^{-\nu\psi}$  with  $\nu \approx 1.2, \psi \approx 0.5$



# Mott transition in a Bose-Einstein condensate



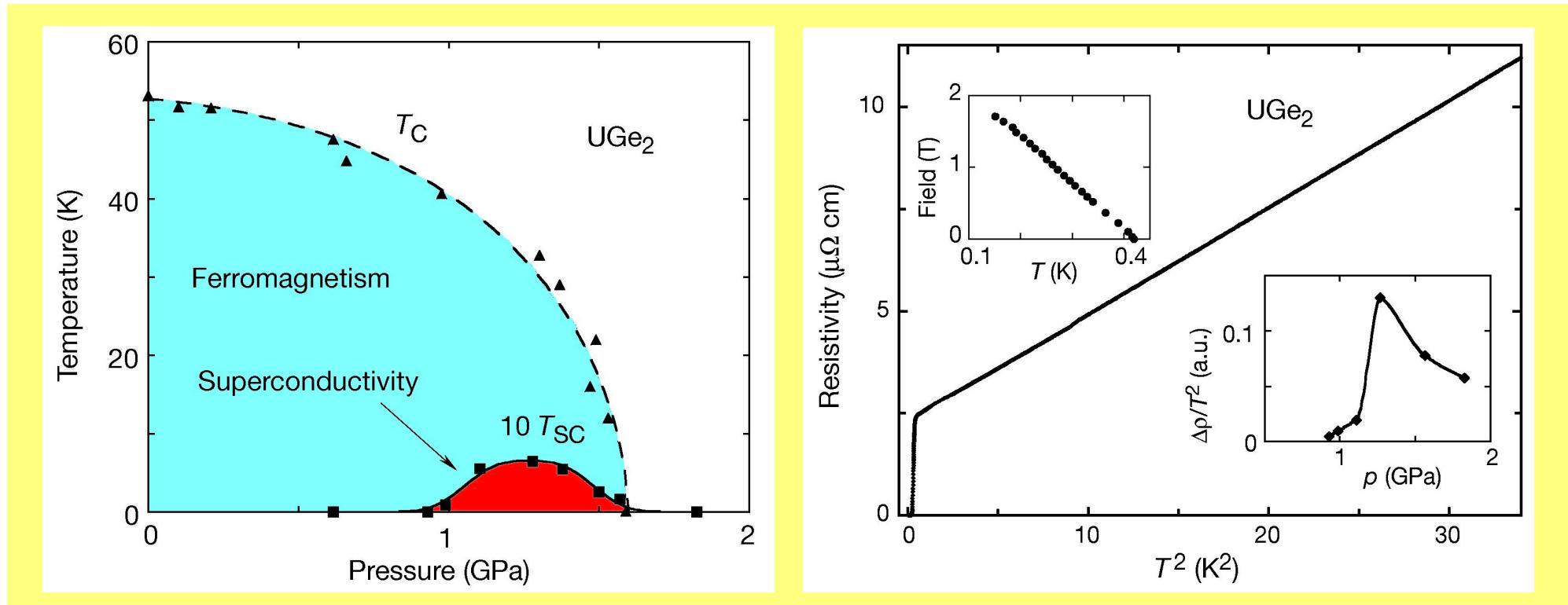
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- Condensed matter physics: complexity and emerging phenomena
  - Phase transitions and quantum phase transitions
  - **Novel phases close to quantum critical points**

# Exotic superconductivity in $\text{UGe}_2$

## Phase diagram:

- phase diagram of  $\text{UGe}_2$  has pocket of **superconductivity** close to ferromagnetic quantum phase transition (electrical resistivity **vanishes** below about 1K)
- in this pocket,  $\text{UGe}_2$  is **ferromagnetic and superconducting** at the same time
- superconductivity appears only in superclean samples



# Character of superconductivity in $\text{UGe}_2$

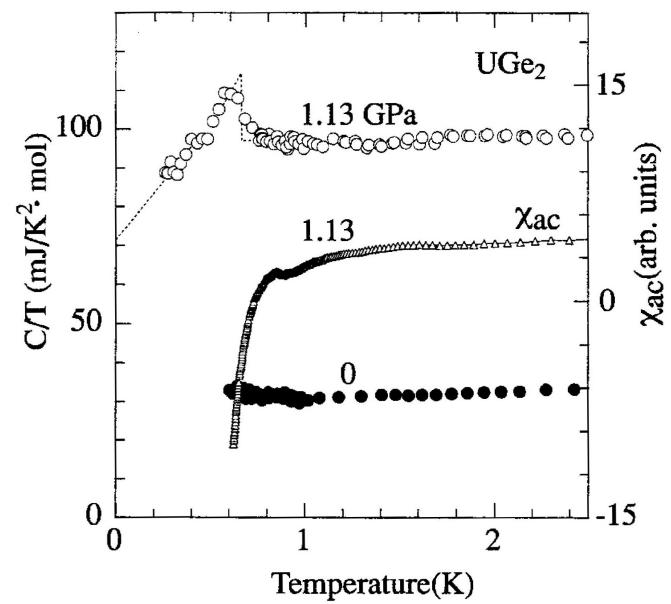
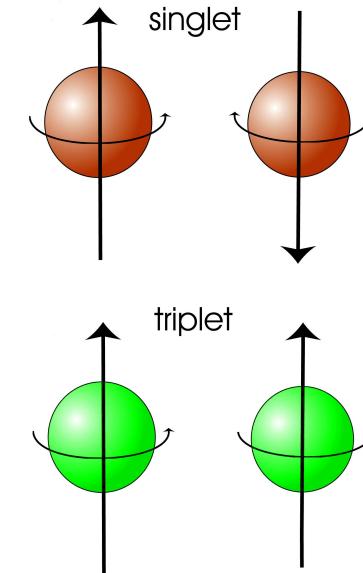
not compatible with conventional (BCS) superconductivity:

- in superconductor, electrons form (Cooper) pairs of spin-up and spin-down electrons
- ferromagnetism requires majority of spins to be in one direction

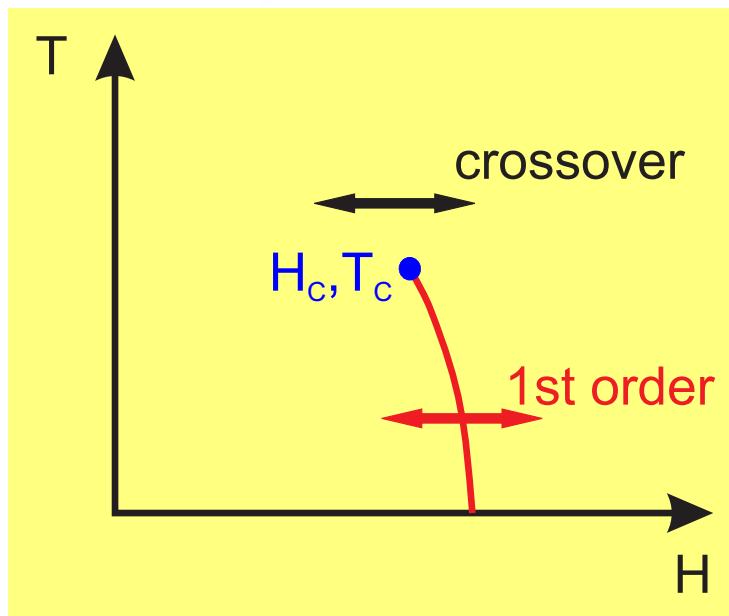
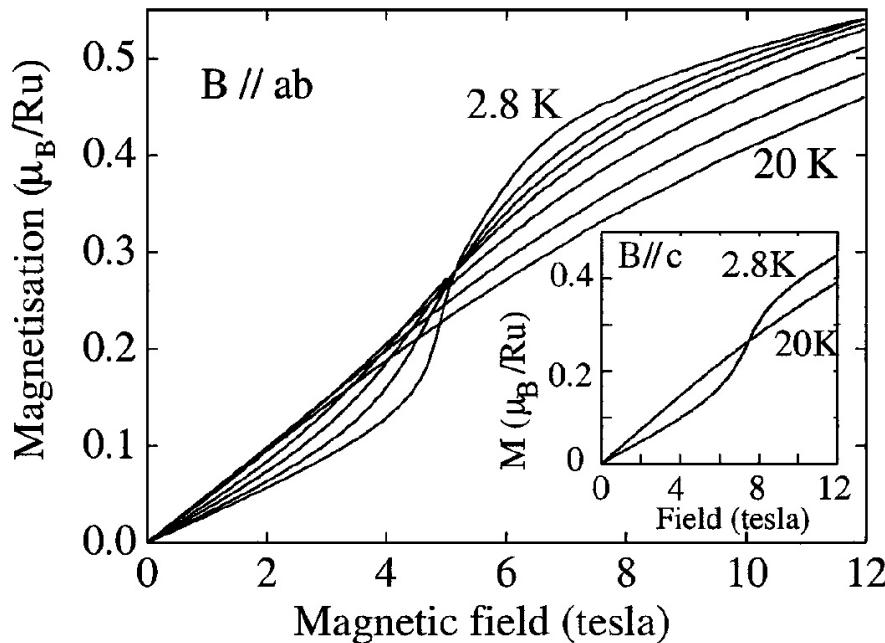
theoretical ideas:

- phase separation (layering or disorder): **NO!**
- partially paired FFLO state: **NO!**
- spin triplet pairs with odd spatial symmetry, magnetic fluctuations promote this type of pairing

Magnetic quantum phase transition induces spin-triplet superconductivity



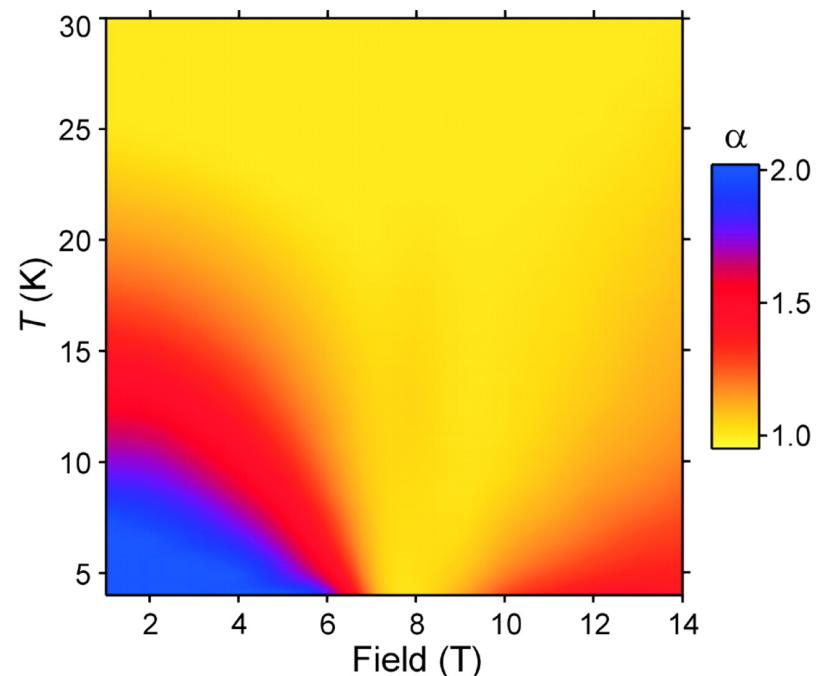
# Metamagnetic transition in $\text{Sr}_3\text{Ru}_2\text{O}_7$



- $\text{Sr}_3\text{Ru}_2\text{O}_7$  undergoes metamagnetic transition as function of field
- critical endpoint can be tuned to  $T = 0$  by tilting the field

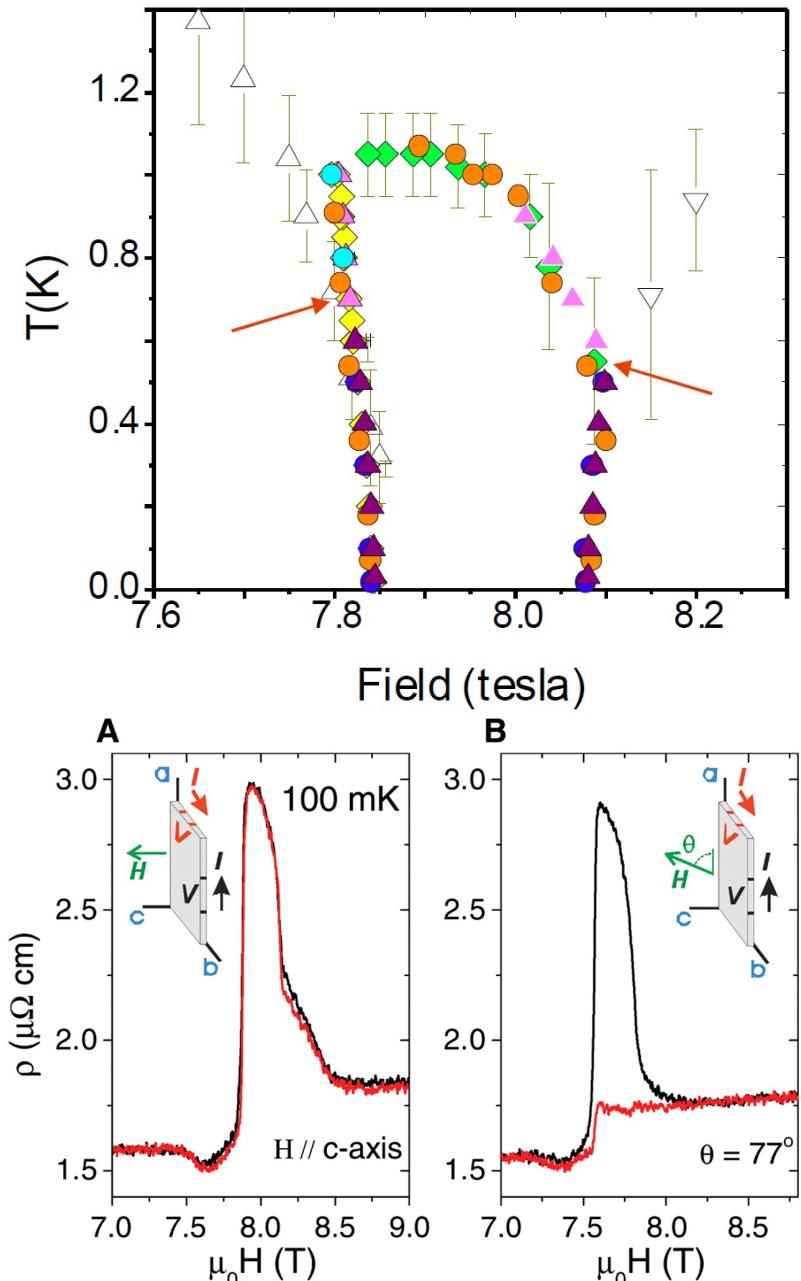
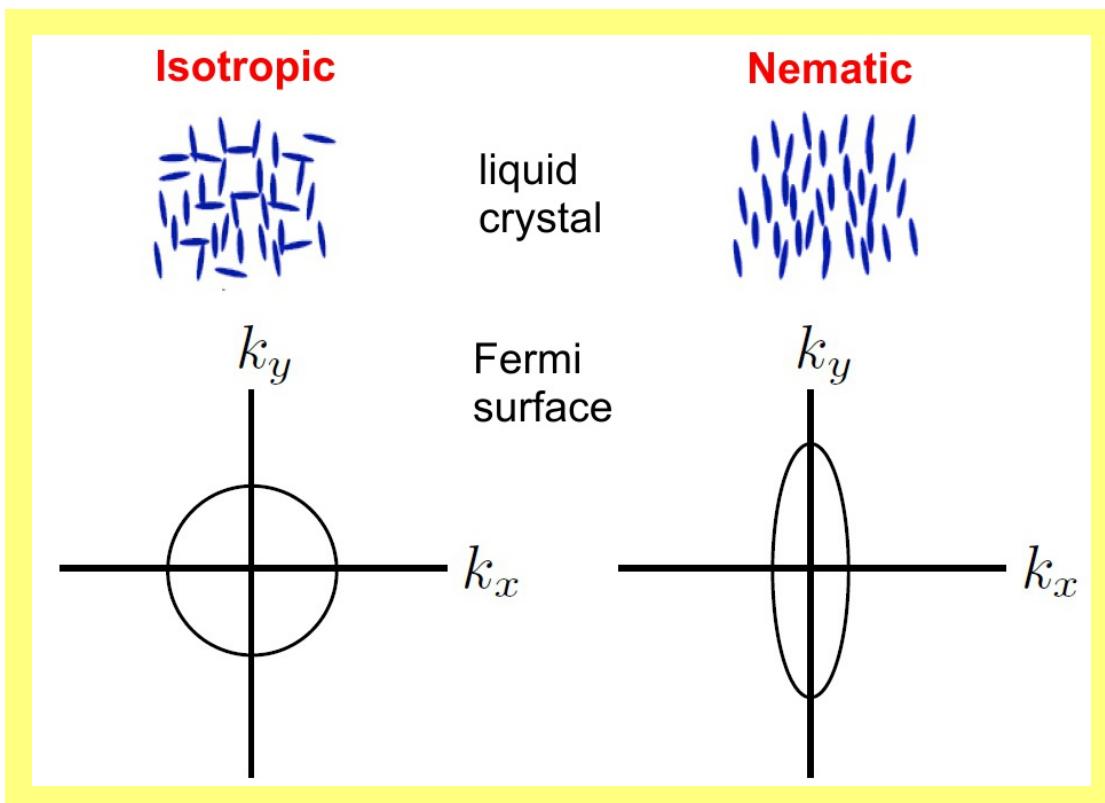
⇒ **metamagnetic quantum critical point**

- seen in temperature dependence of resistivity  $\rho = \rho_0 + AT^\alpha$



# Electronic nematic phase in $\text{Sr}_3\text{Ru}_2\text{O}_7$

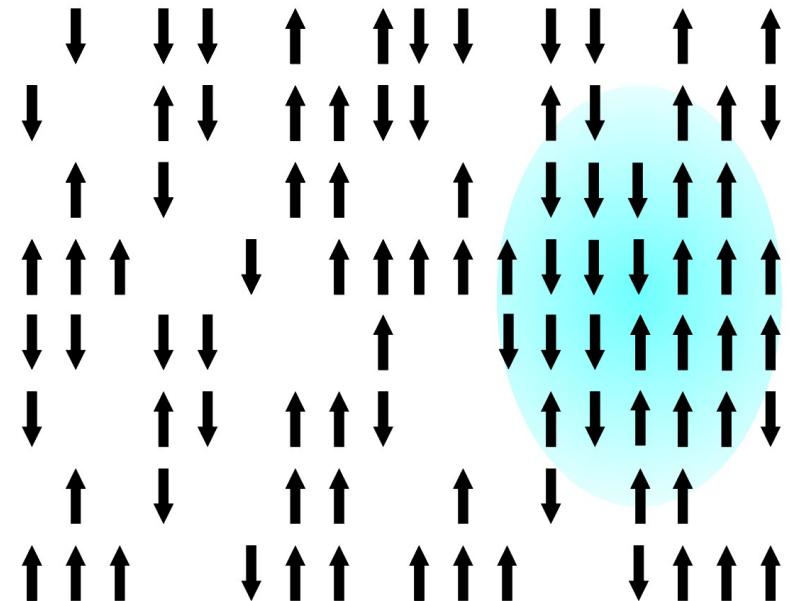
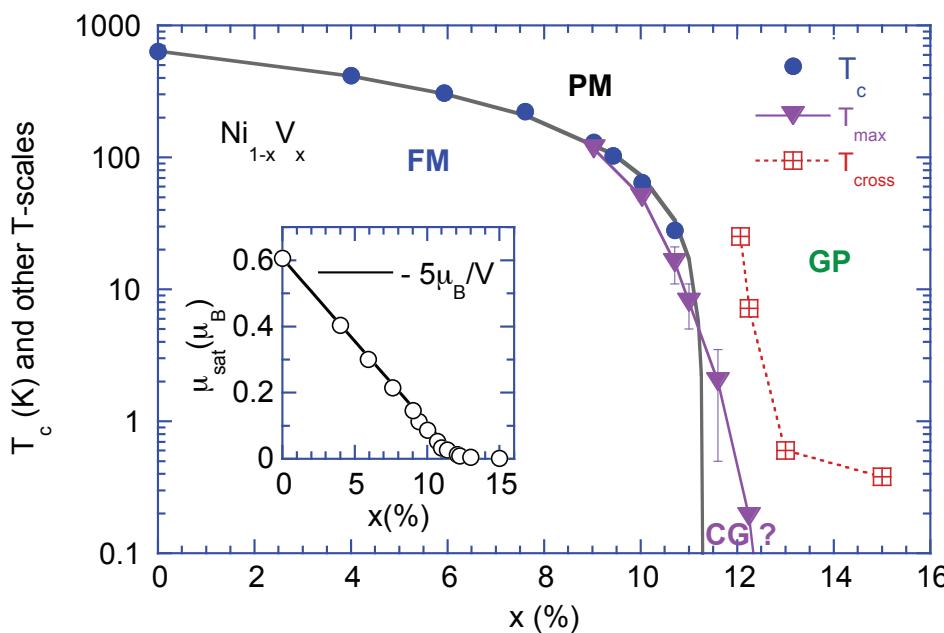
- in pure samples at low temperatures, QCP **preempted by novel phase**
- resistivity highly anisotropic ( $\text{C}4 \rightarrow \text{C}2$ )  
⇒ new phase is **electronic nematic** (translational invariant but rotational symmetry spontaneously broken)



# Disorder and Griffiths phases

## QPT in a disordered system:

- **rare region** can be locally in ordered phase even if bulk system is in disordered phase
- probability of rare region **exponentially small**  $p(L) \sim \exp(-cL^d)$ :
- rare regions act as large **superspins**  
⇒ slow dynamics, large contribution to TD

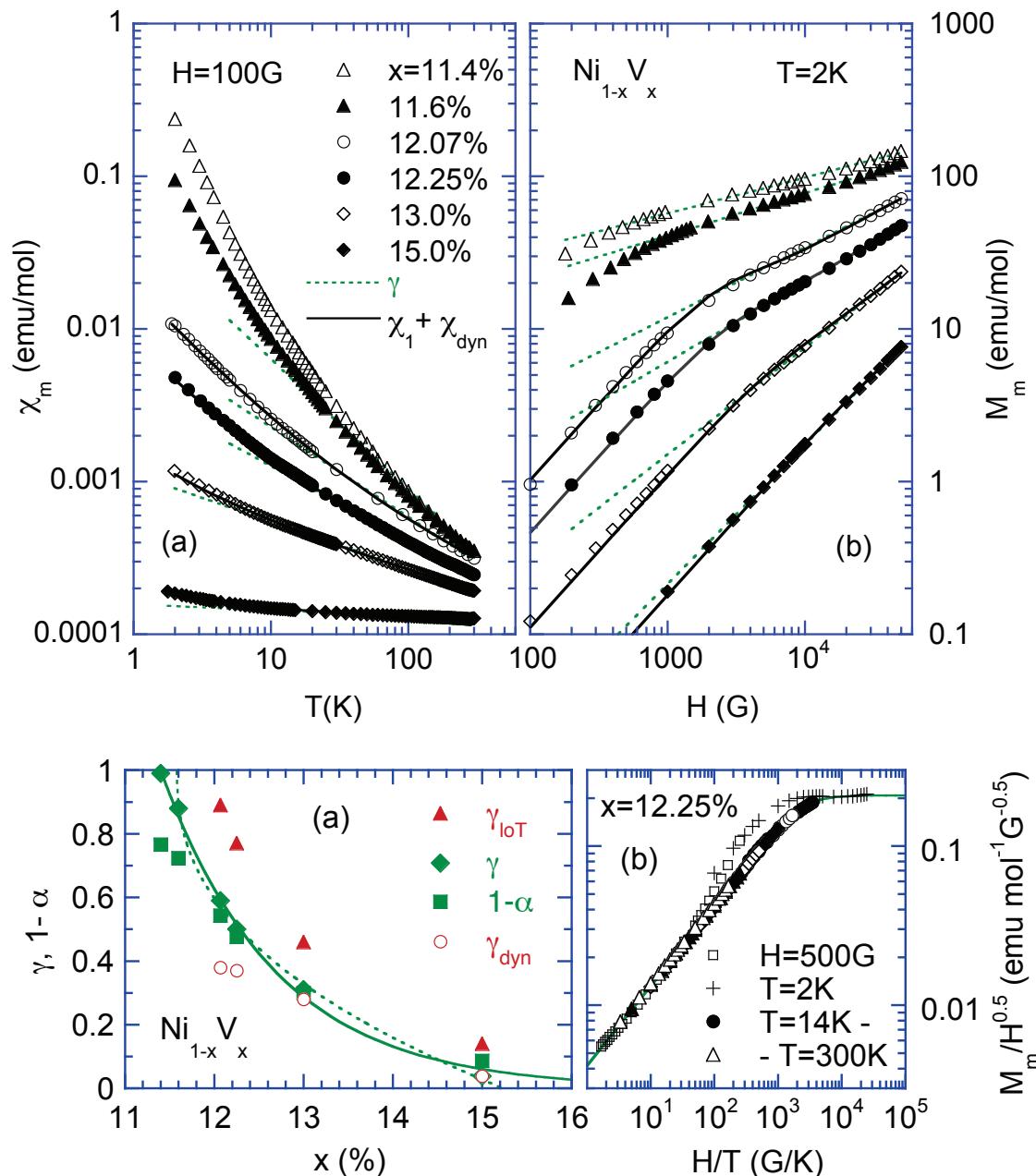


Can rare regions dominate the thermodynamic response?

⇒ **quantum Griffiths phase**

- example:  
diluted ferromagnet  $\text{Ni}_{1-x}\text{V}_x$

# Quantum Griffiths phase in $\text{Ni}_{1-x}\text{V}_x$

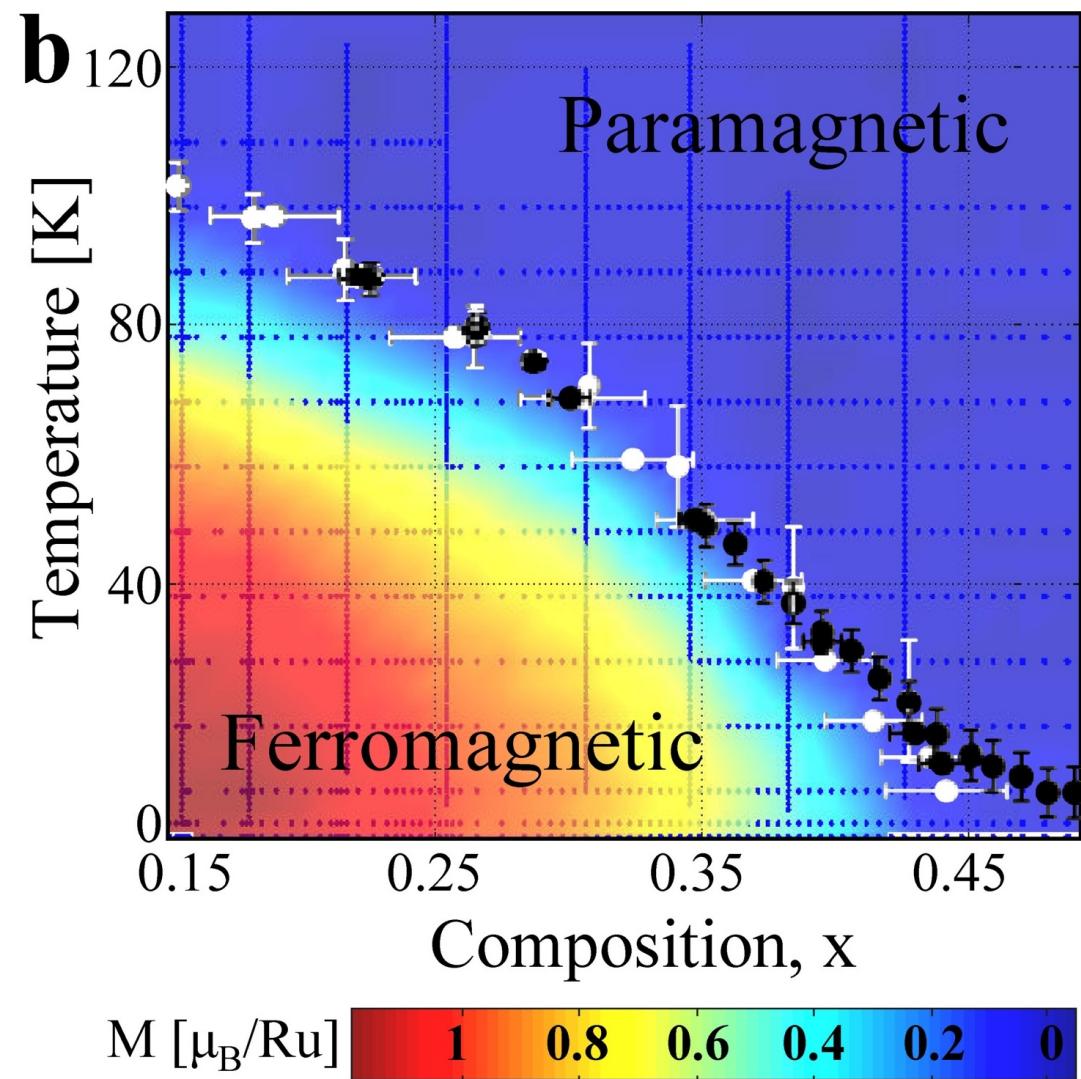
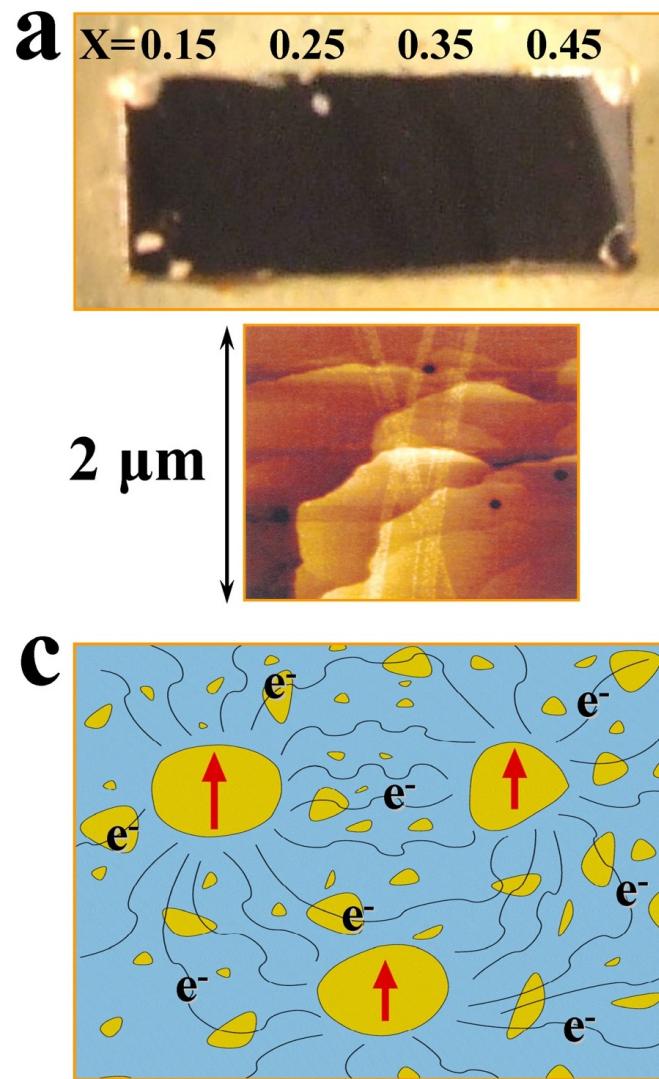


at concentrations above  $x_c$ :

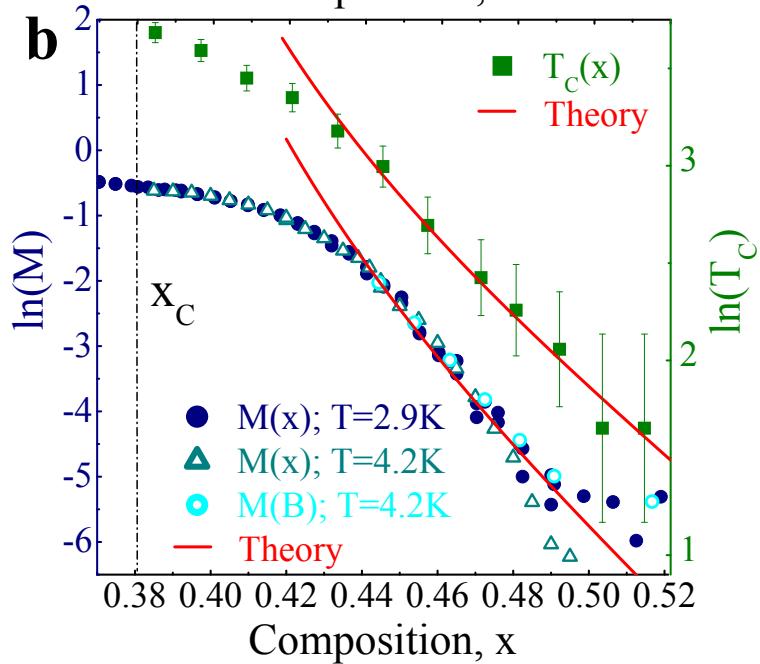
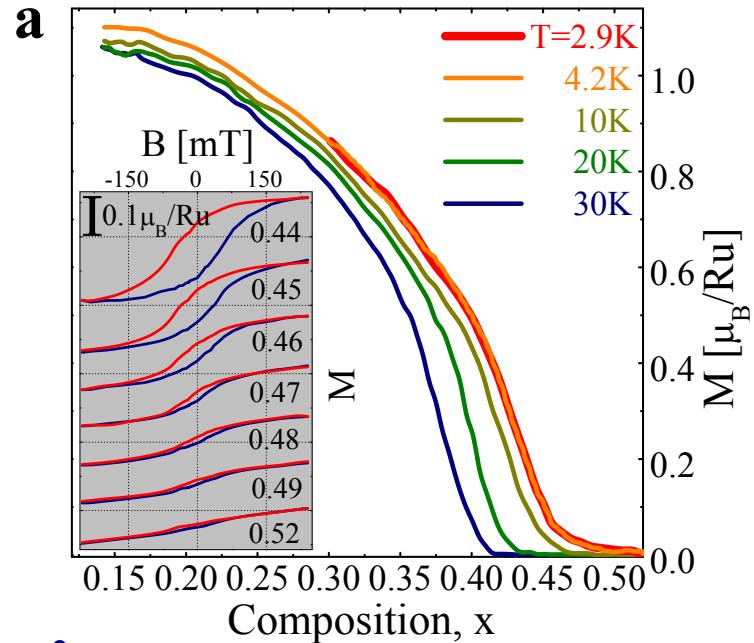
- strongly enhanced magnetic response
- susceptibility  $\chi(T)$  diverges as  $T \rightarrow 0$
- $\chi(T)$  and  $m(H)$  follow nonuniversal power laws  $\chi \sim T^{\lambda-1}$ ,  $m \sim H^\lambda$  (**Griffiths singularities**)
- Griffiths exponent  $\lambda = 1 - \gamma$  varies systematically with  $x$
- experiments agree with **infinite-randomness** critical point scenario

quantum Griffiths phase for  $x \approx 11.5$  to 15%.

# Rare regions and smeared phase transition in $\text{Sr}_{1-x}\text{Ca}_x\text{RuO}_3$



# Composition-tuned smeared phase transitions

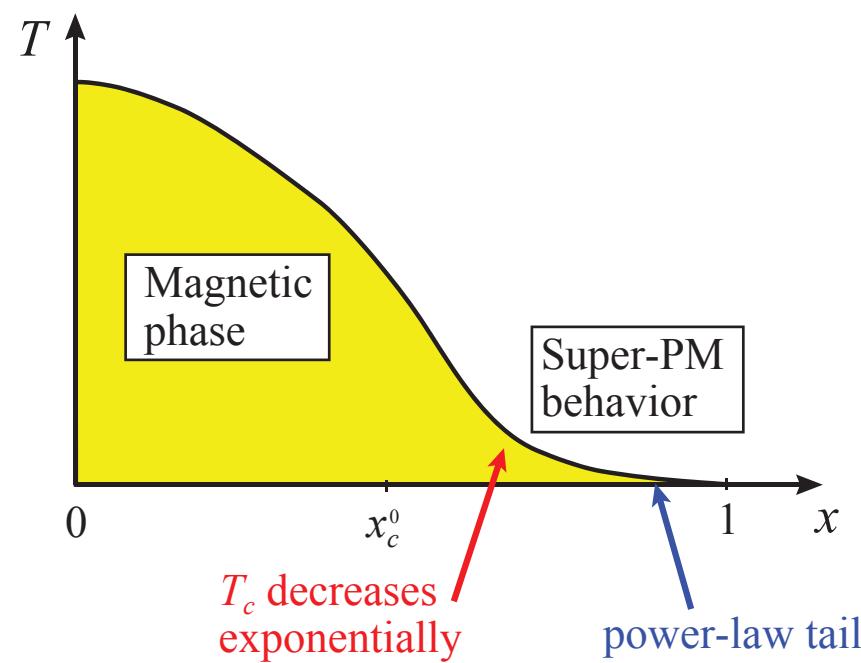


Magnetization and  $T_c$  in tail:

$$M, T_c \sim \exp \left[ -C \frac{(x - x_c^0)^{2-d/\phi}}{x(1-x)} \right]$$

for  $x \rightarrow 1$ :

$$M, T_c \sim (1-x)^{L_{\min}^d}$$

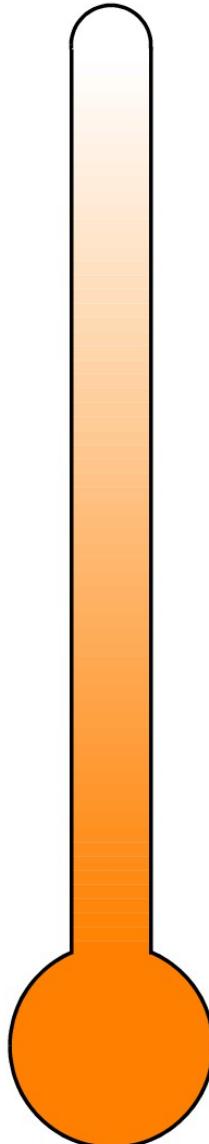


# Conclusions

- “**More is different:**” condensed matter physics explores **emerging phenomena** caused by the interplay of many constituent particles
- new states of quantum matter can be found at **low temperatures** and at **boundaries** between existing phases
- quantum phase transitions occur at **zero temperature** as a function of a parameter like pressure, chemical composition, disorder, magnetic field
- quantum phase transitions are caused by **quantum fluctuations** (i.e, Heisenberg's uncertainty principle) rather than thermal fluctuations
- quantum phase transitions can have fascinating consequences including the genesis of **new phases of matter**

**Quantum phase transitions provide a novel ordering principle in condensed matter physics**

# Wonderland at low temperatures



273K (0C)	water freezes
195K (-78C)	carbon dioxide sublimates (dry ice)
133K (-140C)	<b>superconductivity</b> in cuprate perovskites
77K (-196C)	nitrogen (air) liquefies
66K (-207C)	nitrogen (air) freezes
4.2K (-268.9C)	helium liquefies
2.2K (-270.9C)	helium becomes <b>superfluid</b>
170 nK	<b>Bose-Einstein condensation</b> of rubidium
0K (-273.1C)	<b>absolute zero of temperature</b>

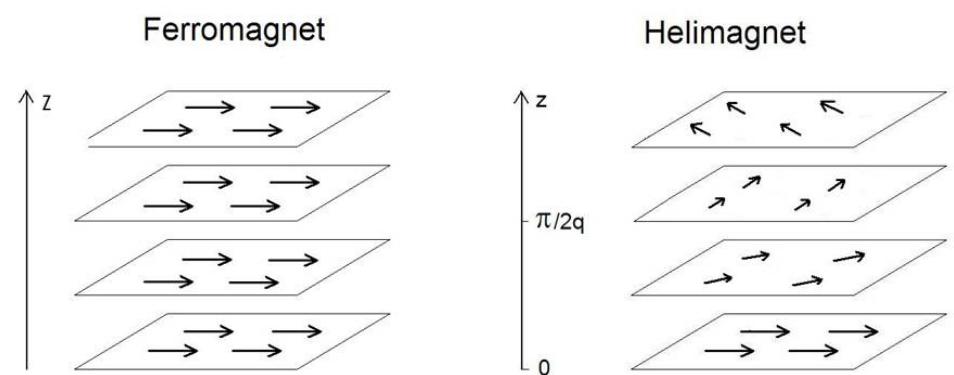
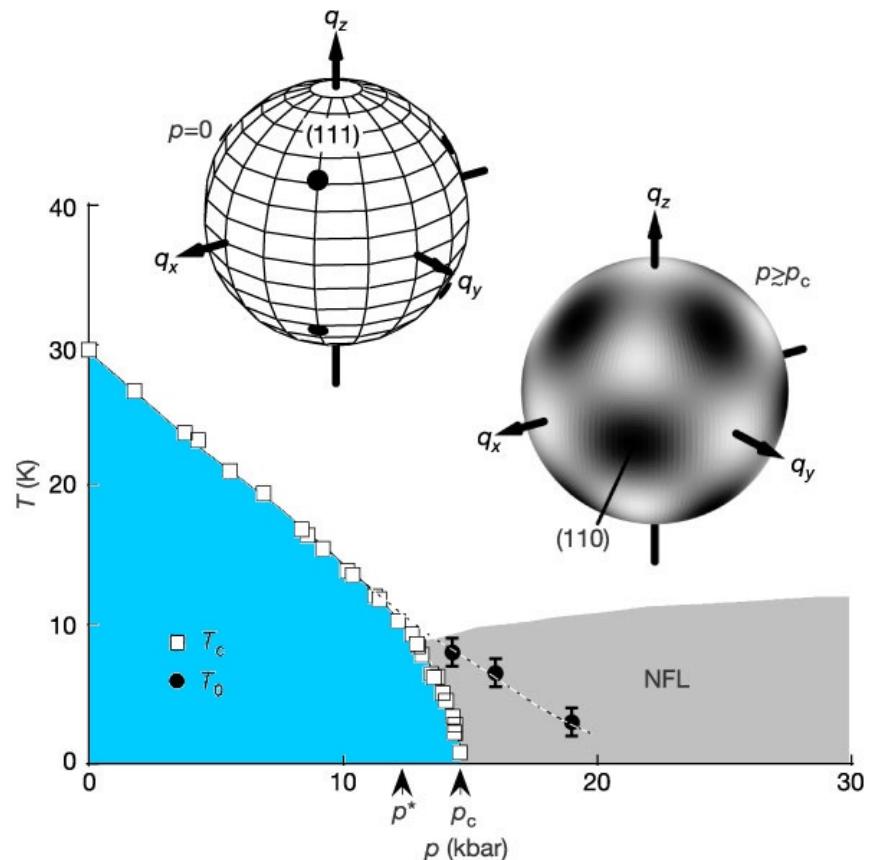
# Magnetic phases in MnSi

## Phase diagram: (Pfleiderer et al, 2004)

- magnetic transition at 30 K at ambient pressure
- transition tunable by hydrostatic pressure
- quantum phase transition at  $p_c = 14$  kbar

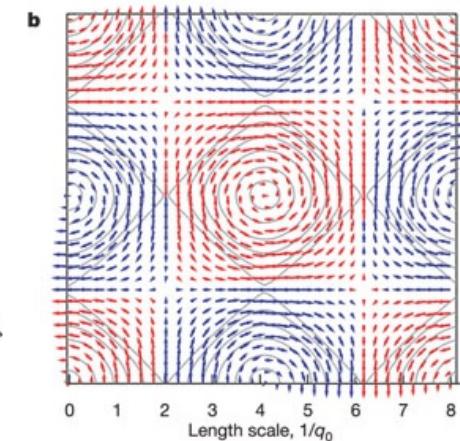
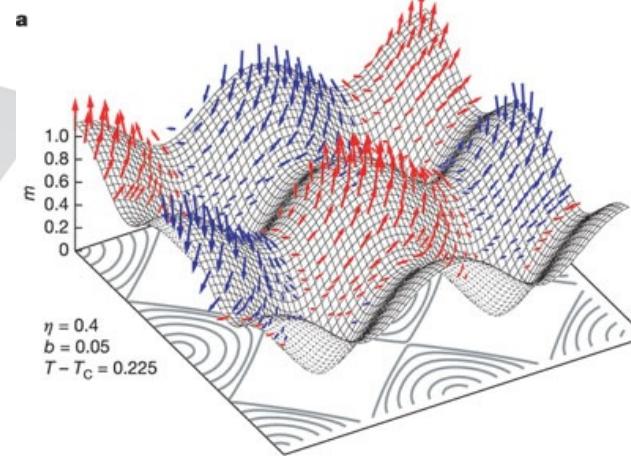
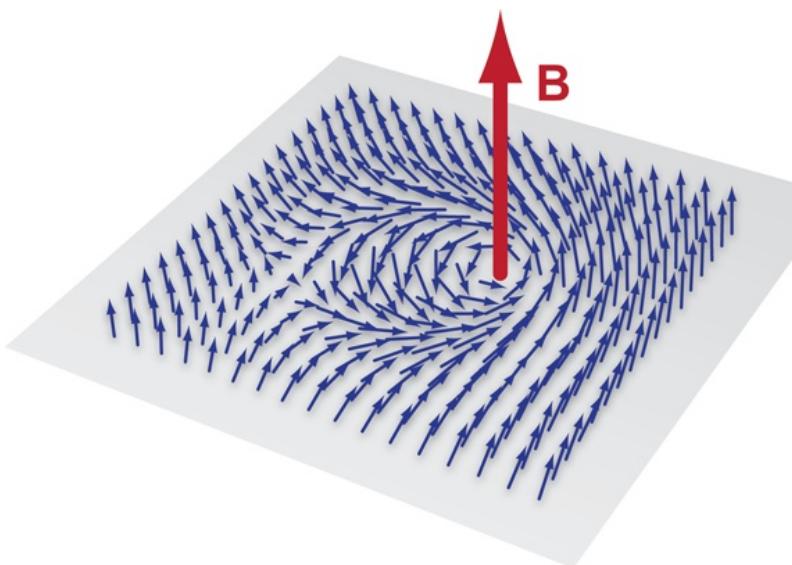
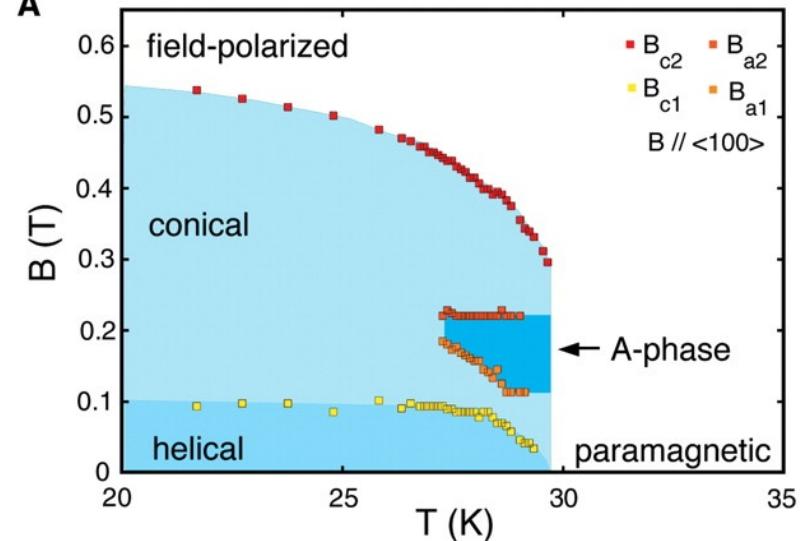
## Magnetic state:

- ordered state is helimagnet with  $q = 180\text{\AA}$ , pinned in (111) direction
- short-range order persists in paramagnetic phase, helical axis depinned



# Skyrmions and skyrmion lattices

- even more exotic magnetic states occur in a magnetic field  $B$
- in “A” phase, magnetization vector forms knots, called **skyrmions**, by twisting in two directions
- these skyrmions arrange themselves into regular **skyrmion lattice**



**If the critical behavior is classical at any nonzero temperature, why are quantum phase transitions more than an academic problem?**

# Phase diagrams close to quantum phase transition

Quantum critical point controls **nonzero-temperature** behavior in its vicinity:

**Path (a):** crossover between classical and quantum critical behavior

**Path (b):** temperature scaling of quantum critical point

